

Applications of Smooth Absorbing Potentials on Mesoscopic Transport

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1 Introduction

- Mesoscopic transport is described by the Caroli formula in terms of Green functions.
- Semi-infinite leads result in energy-dependent self-energies, demanding a new calculation for each energy.
- We propose a two-step approach:
 - 1- Replace the leads by finite buffers with a smooth absorbing potential, yielding energy-independent self-energies
 - 2 - Expand Green functions in a Faber polynomial series, allowing the evaluation of the transmission for many energies simultaneously.

➔ **Result: accurate transmission with a highly efficient, energy-separable algorithm**

2 Model and Caroli Formula

We consider a 1D tight-binding chain with Hamiltonian

$$H = \sum_i \varepsilon_i |i\rangle\langle i| - t \sum_{\langle i,j \rangle} (|i\rangle\langle j| + |j\rangle\langle i|).$$

Transmission is given by the Caroli formula

$$T(E) = \text{Tr} [\Gamma_L(E) G^r(E) \Gamma_R(E) G^a(E)]$$

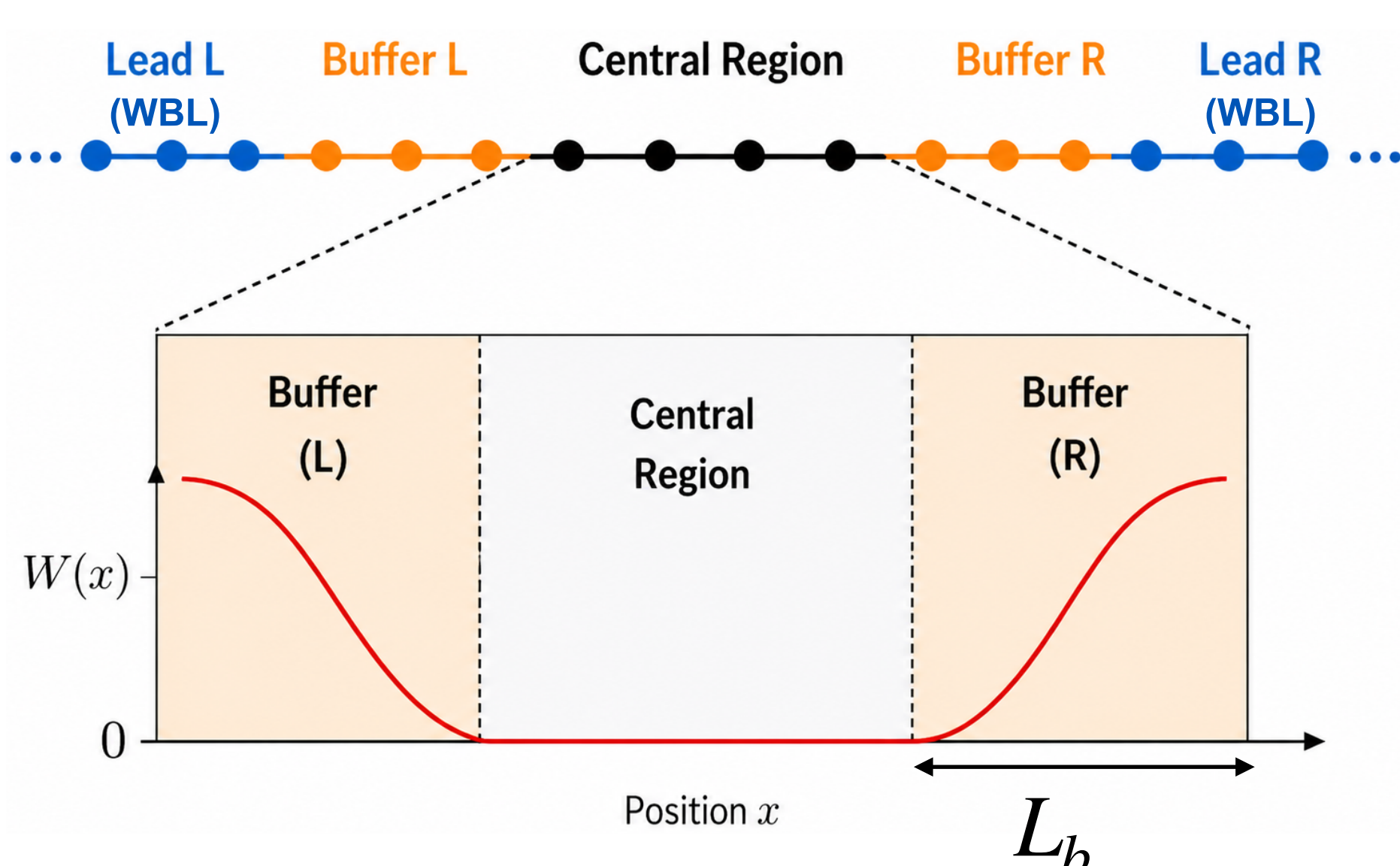
Where

$$- G^r(E) = [EI - H - \Sigma_L(E) - \Sigma_R(E)]^{-1}$$

$$- \Gamma_{L/R}(E) = i [\Sigma'_{L/R}(E) - \Sigma^a_{L/R}(E)]$$

3 Step I - Smooth Absorbing Potential

Semi-infinite leads are replaced by finite buffers with a smooth absorbing potential



The absorbing potential:

$$W(x) = \begin{cases} \eta \left(\frac{x-x_0}{L_b} \right)^n, & x \in \text{left buffer} \\ \eta \left(\frac{x-x_N}{L_b} \right)^n, & x \in \text{right buffer} \end{cases}$$

The self-energies are

$$\Sigma'_{L/R}(x) = -iW_n^{L/R}(x)$$

With sufficiently smooth potential (large L_b) the self-energies are energy independent:

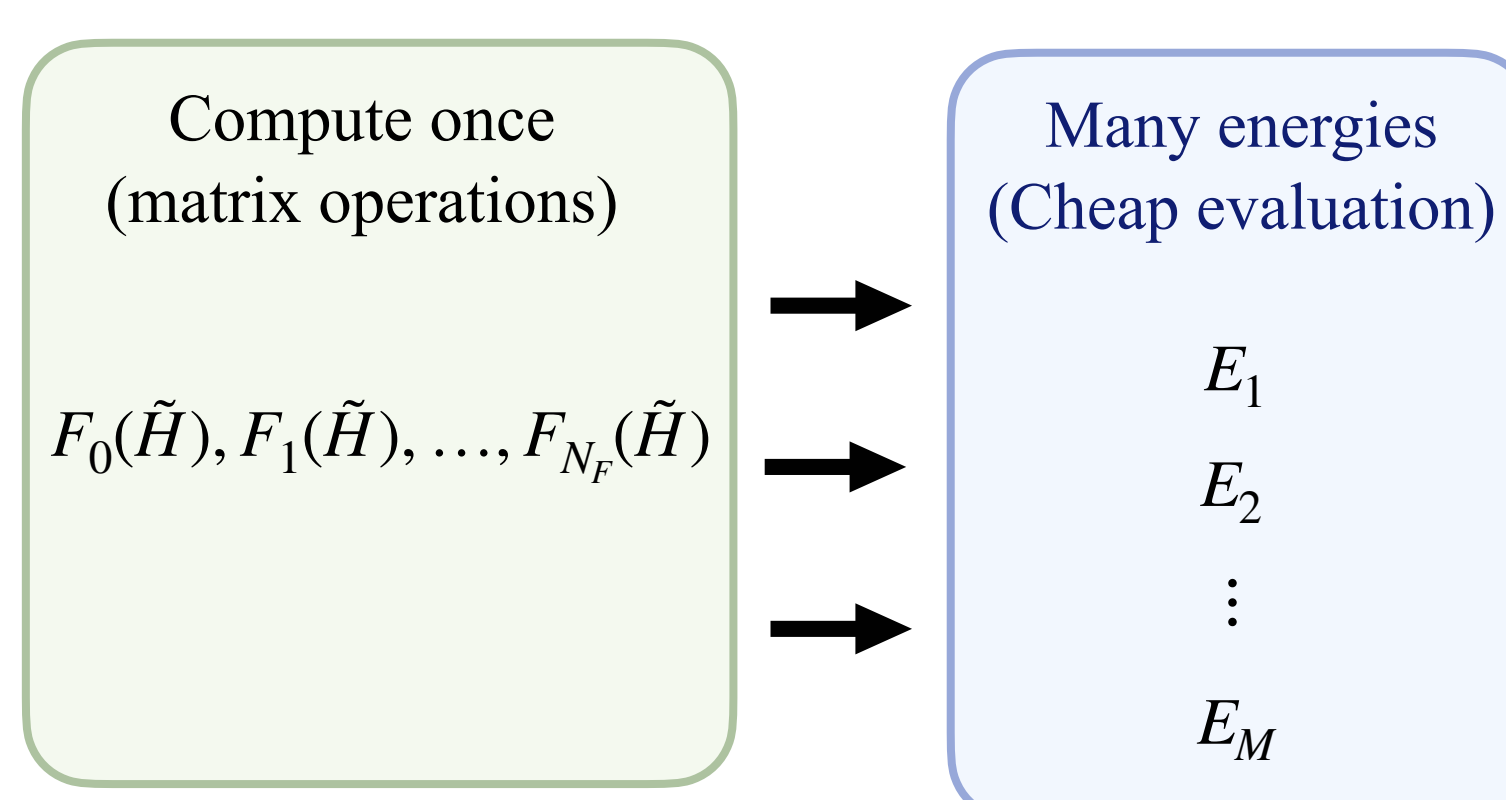
$$\Sigma_{L/R}(E) \approx \Sigma_{L/R}(\text{constant})$$

4 Step II - Faber Expansion

We expand the retarded (advanced) Green function in Faber polynomials

$$G^r(E) \approx \sum_{n=0}^{N_F} g_n(E) F_n(\tilde{H})$$

- Coefficients depend only on energy
- Matrices are computed once and reused for all energies



➔ **One expansion → transmission for many energies simultaneously!**

5 Results

5.1 Step I - Smooth Absorbing Potential

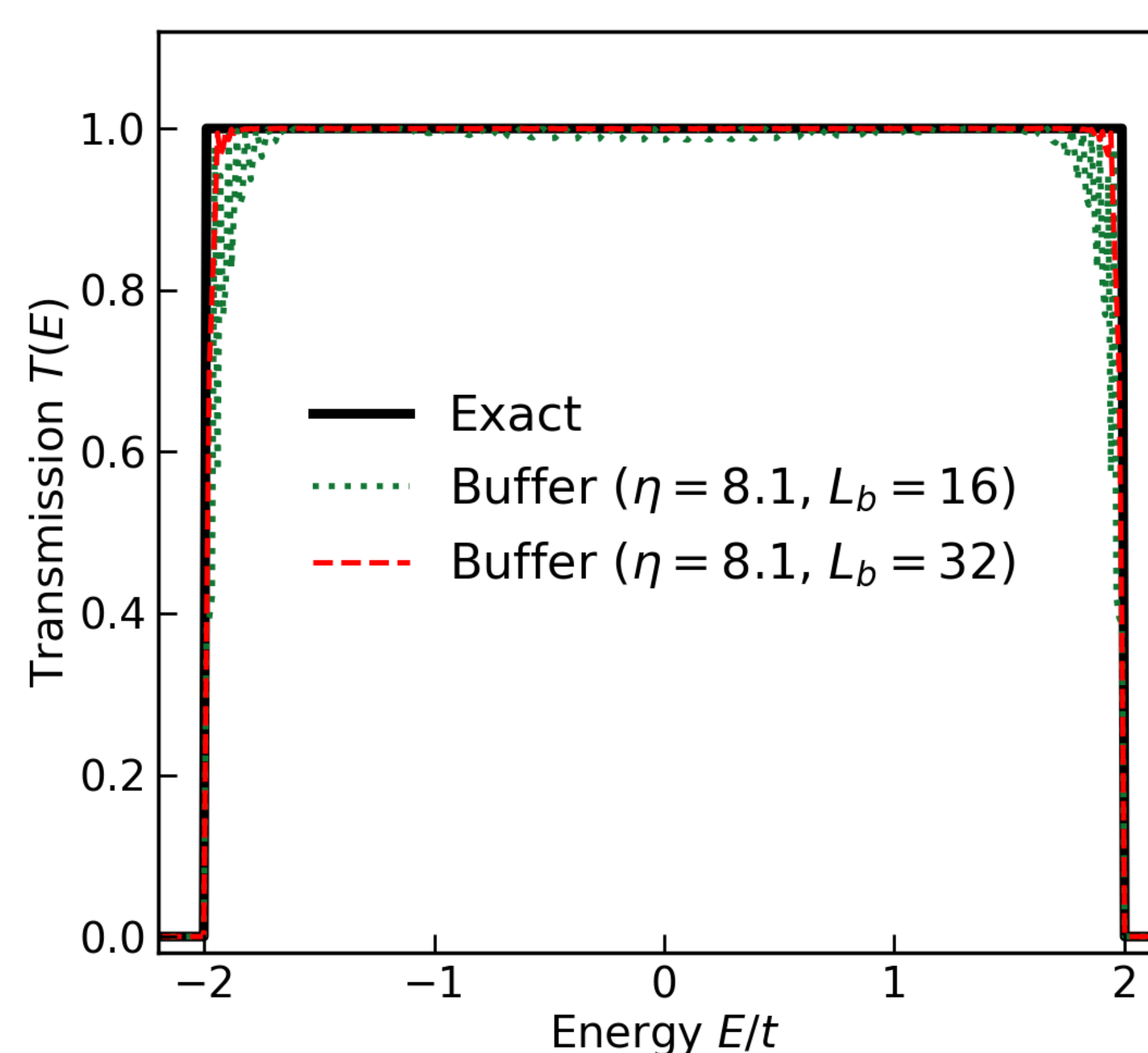


Fig. 2 - Transmission obtained with buffer zoen

The power law absorbing buffer reproduces the exact transmission with good accuracy.

5.2 Step II - Transmission with Faber expansion

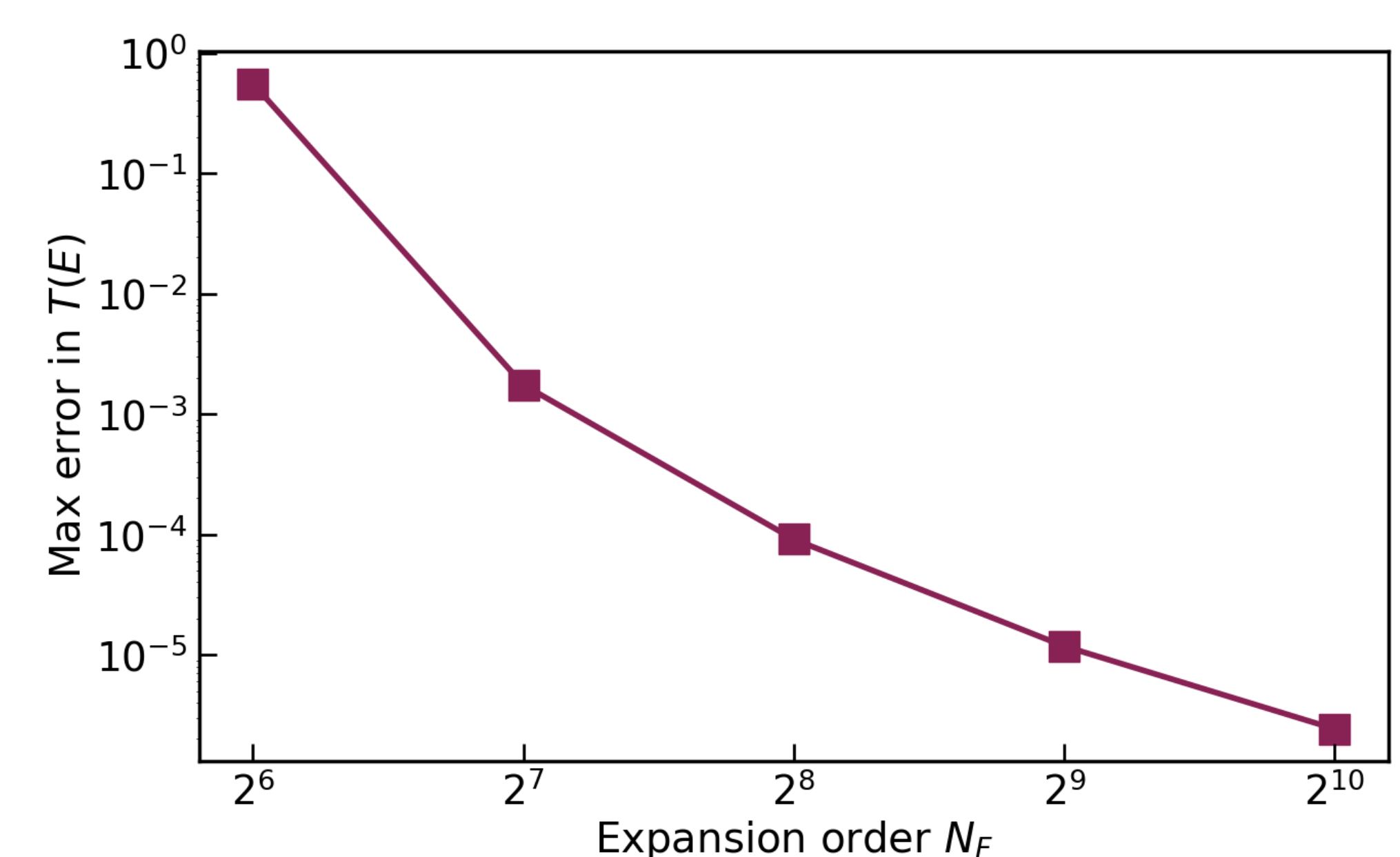
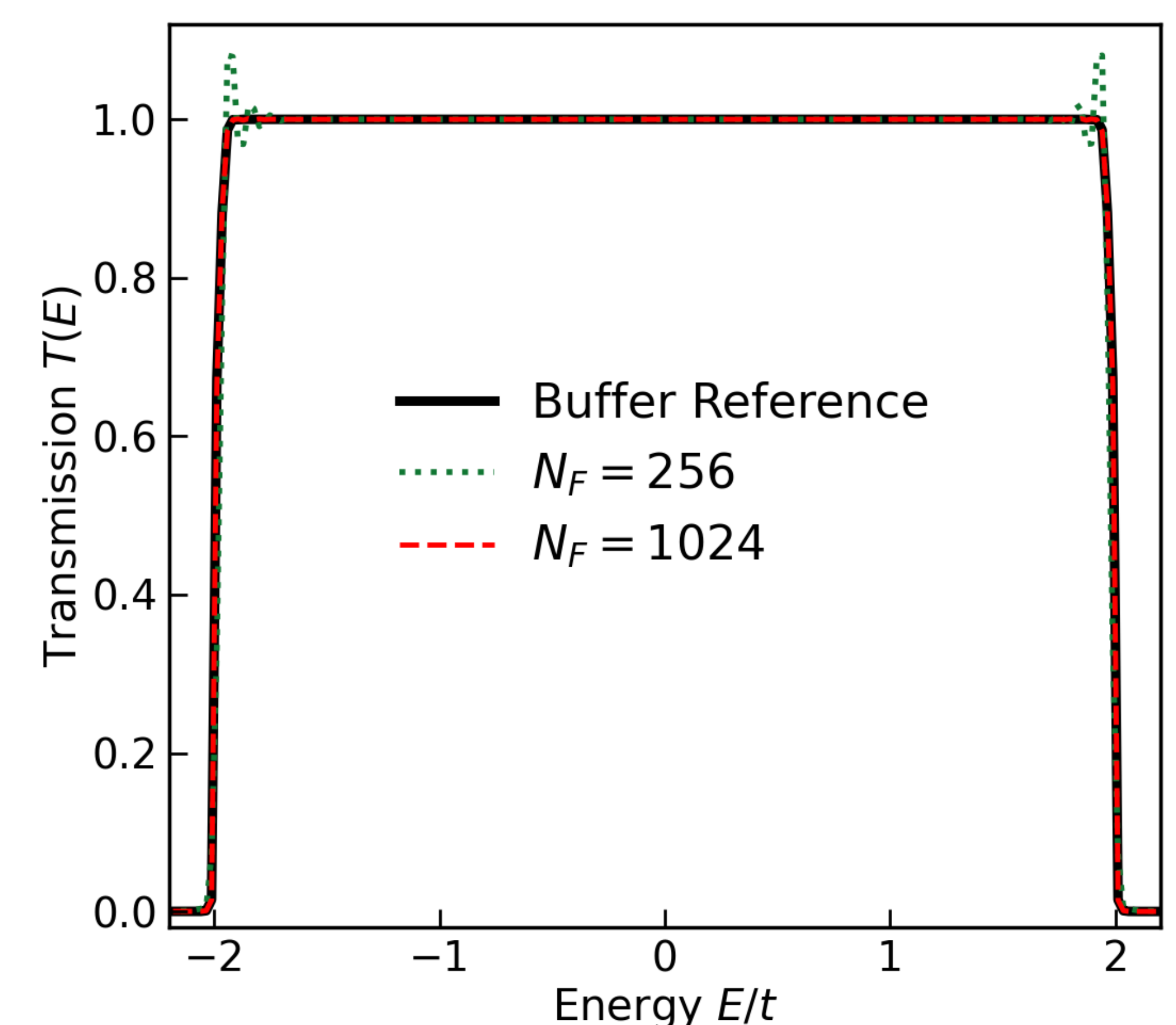


Fig. 4 - **Top** Transmission with finite buffers expanded for different polynomial numbers **Bottom** relative error with expansion order

Faber expansion converges rapidly and accurately reproduces the transmission over

6 Conclusions

- Energy-independent self-energies: no need to recompute $\Sigma_{L/R}(E)$ for each energy;
- Large speedup for dense energy grid;
- Controlled accuracy via buffer size L_b , power order n and expansion order N_F ;

➔ **A scalable and reliable framework for quantum transport calculations**

7 References

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8 Acknowledgements

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