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Modelling realistic open quantum systems with process tensors



Jonathan Keeling
University of St Andrews
Lisbon, April 2026



Funded by
the European Union



QOpen

Many-body Open Quantum Systems

COST Action CA24109 - QOpen

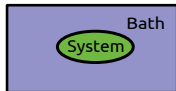
<https://cost.eu>

<https://qopen.eu/>

COST (European Cooperation in Science and Technology) is a funding agency for research and innovation networks. Our Actions help connect research initiatives across Europe and enable scientists to grow their ideas by sharing them with their peers. This boosts their research, career and innovation.

Open quantum systems & Markovian approx.

- Want $\rho_S(t) = \text{Tr}_B(\rho(t))$.



- Standard OQS derivation:

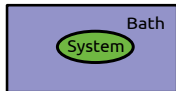
- Time local EOM:

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- $H = H_S + \sum_k X_S g_k (\hat{b}_k^\dagger + \hat{b}_k) + \omega_k \hat{b}_k^\dagger \hat{b}_k$

- Weak coupling (Born approx.)

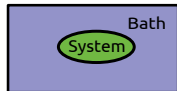
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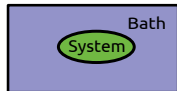
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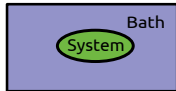
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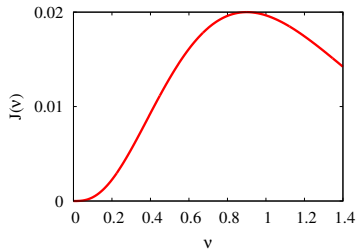
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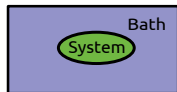


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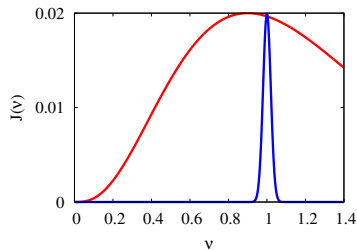
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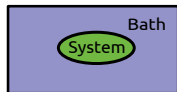
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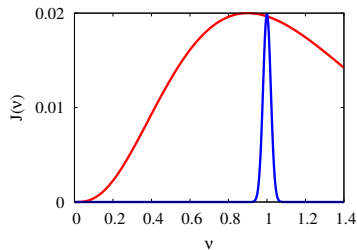
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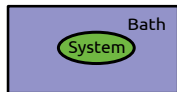


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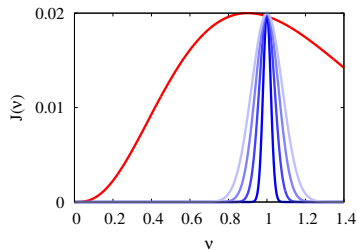


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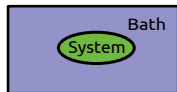


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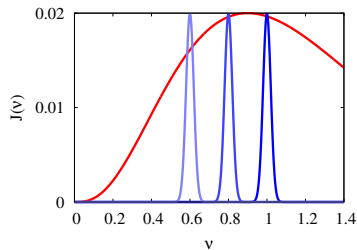


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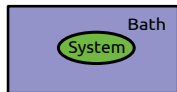


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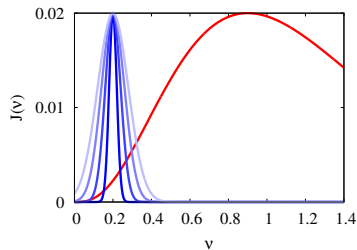


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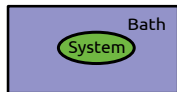


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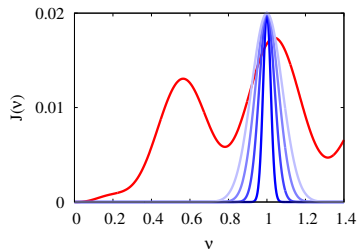
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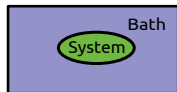
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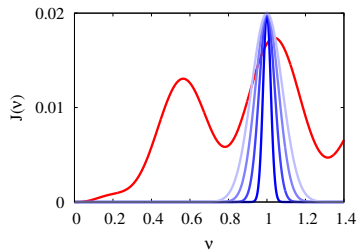
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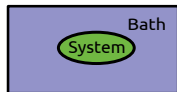
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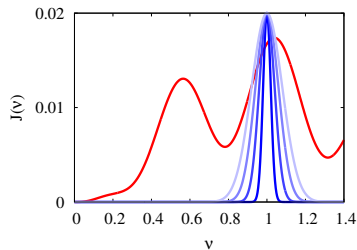
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- Exactly soluble problems ...

- e.g. Bosonic: $H = \Psi_I^\dagger M_{IJ} \Psi_J + \sum_{I,k} \xi_{I,k} (\psi_I + \psi_I^\dagger)(b_k^\dagger + b_k) + H_{\text{bath}}$

- Independent Boson model, $H = \sigma_z(A + \sum_k \xi_k(b_k^\dagger + b_k)) + H_{\text{bath}}$

- Polaron master equation

$$H \rightarrow e^{-V} H e^V, \quad V = \sum_k \xi_k (b_k - b_k^\dagger) X_{\text{sys}}$$

- Renormalize system parameters

- Perturbative remaining coupling

[Jang, J. Chem. Phys '09, McCutcheon *et al.* PRB '11, Roy and Hughes PRB '12]

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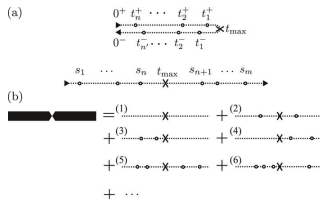
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non-Markovian methods: augmentation

- Exponential bath correlations, Hierarchical EOM [Tanimura & Kubo, JPSJ '89]
Bath $\rightarrow$ sum of Lorentzians modes. n_k describes interaction with mode k .

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- Hierarchy of Pure States [Suess, Eisfeld & Strunz, PRL '14; J. Stat. Phys. '15]

Stochastic unravelling $\rho_{\mathbf{n},\mathbf{m}} \rightarrow \sum_i |\psi_i^{\mathbf{n},\mathbf{m}}\rangle \langle \psi_i^{\mathbf{n},\mathbf{m}}|$.

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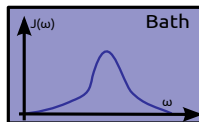
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System



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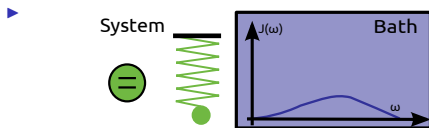
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► Include bath $\rightarrow$ chain mapping
[Chin *et al.* PRL '10; Nat. Phys. '13; Woods *et al.*, J. Math. Phys. '14]

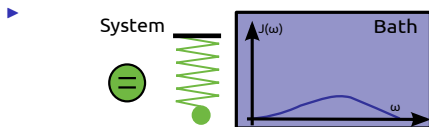
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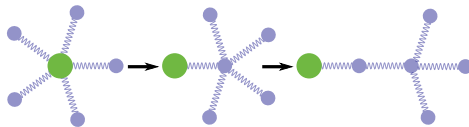
$$\dot{\rho}_{\mathbf{n},\mathbf{m}} = \mathcal{L}_S \rho_{\mathbf{n},\mathbf{m}} + \sum_{\mathbf{n}',\mathbf{m}'} \mathcal{D}_{\mathbf{n},\mathbf{m},\mathbf{n}',\mathbf{m}'} \rho_{\mathbf{n}',\mathbf{m}'}$$

- Hierarchy of Pure States [Suess, Eisfeld & Strunz, PRL '14; J. Stat. Phys. '15]
Stochastic unravelling $\rho_{\mathbf{n},\mathbf{m}} \rightarrow \sum_t |\psi_{\mathbf{n}}^t\rangle \langle \psi_{\mathbf{m}}^t|$.
- Increase system — include resonant modes



[Garraway PRA '97, Iles-Smith *et al.* PRA '14, Schröder *et al.* '17]

- Include bath $\rightarrow$ chain mapping
[Chin *et al.* PRL '10; Nat. Phys. '13; Woods *et al.*, J. Math. Phys. '14]



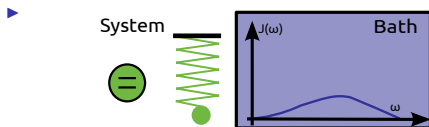
• Augment state space/history: QUAPI [Maki and Makarov, J. Chem Phys '95]

non-Markovian methods: augmentation

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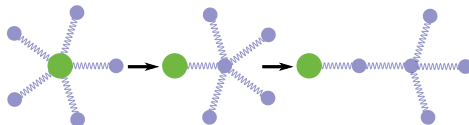
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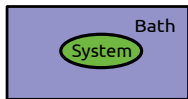


- Augment state space/history: QUAPI [Makri and Makarov, J. Chem Phys '95]

- 1 Non-Markovian problems & Many-body physics
 - Non-Markovian methods
- 2 Background: QUAPI & Tensor networks
 - Path integral approach
 - Tensor networks & Matrix product states
 - The TEMPO algorithm
- 3 Re-interpretation: Process tensors
 - PT-MPO representation
 - Beyond Gaussian environments
- 4 Applications
 - Optimisation
 - Lattice models
 - Polariton condensation
 - Two-dimensional spectroscopy
- 5 Future outlook

QUAPI

- System + Harmonic bath



$$H = H_S + \sum_k g_k X_S (b_k + b_k^\dagger) + \omega_k b_k^\dagger b_k$$

- Write $\rho(t)$ as path sum/integral

- Discrete sum for system

$$1 = \sum_j |j\rangle\langle j|$$

- Path integral for bath

$$1 = \int dx |x\rangle\langle x|$$

- Discretize times, $t_N = N\Delta$

- Integrate out bath.

[Marki & Makarov J. Chem. Phys. '95]

- Path sum (doubled indices) $j = (j_r, j_b)$:

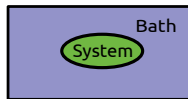
$$\rho_{IN}(t_N) = \sum_{j_0 \dots j_{N-1}} \left(\prod_{n=1}^N \prod_{k=0}^{n-1} I_k(j_n, j_{n-k}) \right) \rho_H(t_0)$$

- Influence functionals:

$$I_k(j, l) = \begin{cases} e^{\phi_k(j, l)} & k \neq 1 \\ e^{\phi_1(j, l)} [e^{\Delta C_0}]_{ll} & k = 1 \end{cases}$$

QUAPI

- System + Harmonic bath



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$$\mathbb{1} = \sum_j |j\rangle\langle j|$$

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[Marki & Makarov J. Chem. Phys. '95]

• Path sum (doubled indices) $J = (j_1, j_0)$:

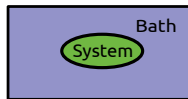
$$\rho_W(t) = \sum_{j_1, \dots, j_{N-1}} \left(\prod_{n=1}^N \prod_{k=0}^{n-1} I_k(j_n, j_{n-k}) \right) \rho_H(t).$$

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QUAPI

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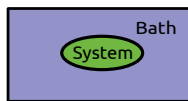
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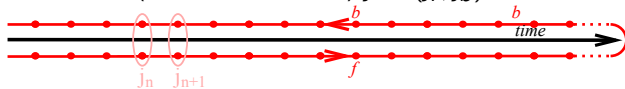
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[Marki & Makarov J. Chem. Phys. '95]

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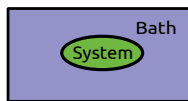
$$\rho_{j_N}(t_N) = \sum_{j_1 \dots j_{N-1}} \left(\prod_{n=1}^N \prod_{k=0}^{n-1} I_k(j_n, j_{n-k}) \right) \rho_{j_1}(t_1).$$

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$$I_k(U, D) = \begin{cases} e^{i\phi_k(U, D)} & k \neq 1 \\ e^{i\phi_k(U, D)} [e^{i\phi_k(U, D)}]_D & k = 1 \end{cases}$$

QUAPI

- System + Harmonic bath



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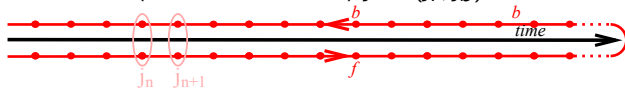
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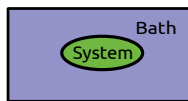
$$I_k(j, j') = \begin{cases} e^{\phi_k(j, j')} & k \neq 1 \\ e^{\phi_1(j, j')} [e^{\Delta \mathcal{L}_0}]_{jj'} & k = 1 \end{cases}$$

• System propagator: $\mathcal{L}_0$

• $\phi_k(j, j')$ depends on $X_j = \langle j | X_S | j \rangle$ and

QUAPI

- System + Harmonic bath



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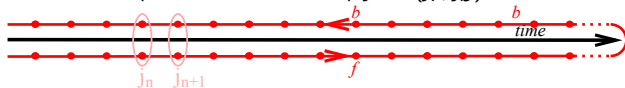
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- $\phi_k(j, j')$ depends on $X_j = \langle j_f | \hat{X}_S | j_b \rangle$ and

$$C(t) = \int d\omega J(\omega) \left[\coth\left(\frac{\omega}{2T}\right) \cos(\omega t) - i \sin(\omega t) \right]$$

QUAPI as tensor propagation

- Evaluating path sum:

$$\rho_{j_N}(t_N) = \sum_{j_1 \dots j_{N-1}} \left(\prod_{n=1}^N \prod_{k=0}^{n-1} I_k(j_n, j_{n-k}) \right) \rho_{j_1}(t_1).$$

- Write as tensor propagation:

$$\rho_{j_N}(t_N) = \sum_{j_1 \dots j_{N-1}} A^{j_N, j_{N-1}, \dots, j_1}$$

- Products $\rightarrow$ Growth.

$$A^{b,j_1} = B^{b,j_1} A^h, \quad A^{b_1,b_2,h} = B^{b_1,b_2,h} A^{b,h}, \quad A^{b_1,b_2,b_3,h} = B^{b_1,b_2,b_3,h} A^{b,b,h}$$

$$B^{b,j_{n-1}, \dots, j_1}_{b_{n-1}, \dots, b_1} = \left(\prod_{k=1}^{n-1} \delta_{j_{n-k}}^{b_{n-k}} \right) \prod_{k=0}^{n-1} I_k(j_n, j_{n-k}).$$

QUAPI as tensor propagation

- Evaluating path sum:

$$\rho_{j_N}(t_N) = \sum_{j_1 \dots j_{N-1}} \left(\prod_{n=1}^N \prod_{k=0}^{n-1} l_k(j_n, j_{n-k}) \right) \rho_{j_1}(t_1).$$

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$$B^{b,j_{n-1}, \dots, j_1} = \left(\prod_{k=1}^{n-1} d_{j_{n-k}}^{j_{n-k-1}} \right) \prod_{k=0}^{n-1} l_k(j_n, j_{n-k}).$$

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- Products $\rightarrow$ Growth.

$$A^{j_2, j_1} = B^{j_2, j_1}_{i_1} A^{i_1}, \quad A^{j_3, j_2, j_1} = B^{j_3, j_2, j_1}_{i_2, i_1} A^{i_2, i_1}, \quad A^{j_4, j_3, j_2, j_1} = B^{j_4, j_3, j_2, j_1}_{i_3, i_2, i_1} A^{i_3, i_2, i_1}$$

$$B^{j_n, j_{n-1}, \dots, j_1}_{i_{n-1}, \dots, i_1} = \left(\prod_{k=1}^{n-1} \delta_{i_{n-k}}^{j_{n-k}} \right) \prod_{k=0}^{n-1} l_k(j_n, j_{n-k}).$$

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- Products $\rightarrow$ Growth. **Problem: n^{th} Tensor size $\sim d^{2n}$**

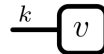
$$A^{j_2, j_1} = B^{j_2, j_1}_{i_1} A^{i_1}, \quad A^{j_3, j_2, j_1} = B^{j_3, j_2, j_1}_{i_2, i_1} A^{i_2, i_1}, \quad A^{j_4, j_3, j_2, j_1} = B^{j_4, j_3, j_2, j_1}_{i_3, i_2, i_1} A^{i_3, i_2, i_1}$$

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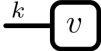
Penrose tensor network notation & contraction

vector

v_k



Penrose tensor network notation & contraction

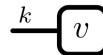
vector v_k 

matrix $A_{i,j}$ 

Penrose tensor network notation & contraction

vector

v_k



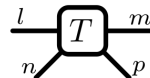
matrix

$A_{i,j}$



4-rank tensor

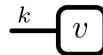
$T_{l,m,n,p}$



Penrose tensor network notation & contraction

vector

v_k



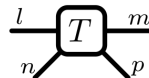
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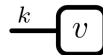
A



Penrose tensor network notation & contraction

vector

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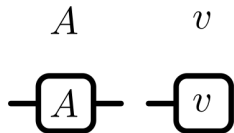
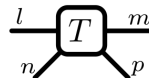
matrix

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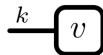


$A_{i,j}$ v_k

Penrose tensor network notation & contraction

vector

v_k



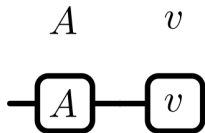
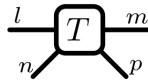
matrix

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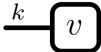
4-rank tensor

$T_{l,m,n,p}$

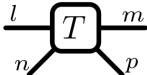


$$\sum_k A_{i,k} v_k$$

Penrose tensor network notation & contraction

vector v_k 

matrix $A_{i,j}$ 

4-rank tensor $T_{l,m,n,p}$ 

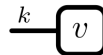
$$A \quad v = w$$


$$\sum_k A_{i,k} v_k = w_i$$

Penrose tensor network notation & contraction

vector

v_k



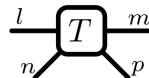
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4-rank tensor

$T_{l,m,n,p}$



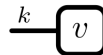
=



Penrose tensor network notation & contraction

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v_k



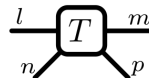
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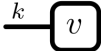


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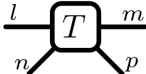
$T_{l,m,n,p}$

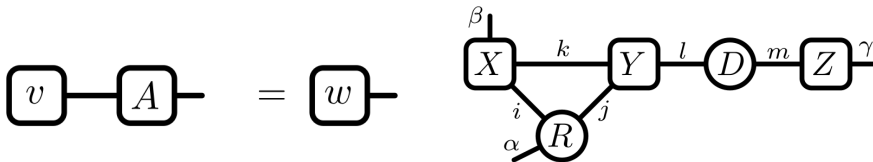


Penrose tensor network notation & contraction

vector v_k 

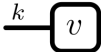
matrix $A_{i,j}$ 

4-rank tensor $T_{l,m,n,p}$ 

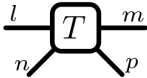


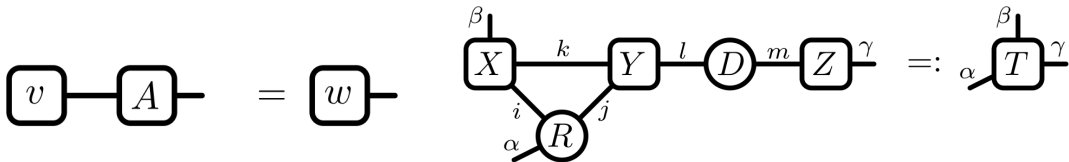
$$\sum_{ijklm} R_{\alpha ij} X_{i\beta k} Y_{jkl} D_{lm} Z_{m\gamma}$$

Penrose tensor network notation & contraction

vector v_k 

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4-rank tensor $T_{l,m,n,p}$ 



$$\sum_{ijklm} R_{\alpha ij} X_{i\beta k} Y_{jkl} D_{lm} Z_{m\gamma} =: T_{\alpha\beta\gamma}$$

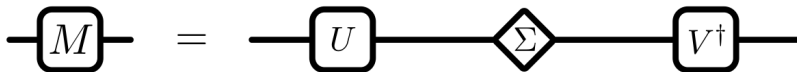
Singular value decomposition & compression



$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}$$

- Fixed bond dimension: Keep $\lambda_1, \dots, \lambda_\chi$
- Fixed precision: Keep $\lambda_n > \epsilon_{\text{rel}} \lambda_1$

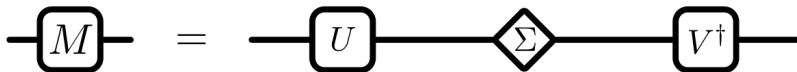
Singular value decomposition & compression



$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} = \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \begin{pmatrix} \lambda_1 & & & \\ & \lambda_2 & & \\ & & \lambda_3 & \\ & & & \lambda_4 \end{pmatrix} \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}$$

- Fixed bond dimension: Keep $\lambda_1 \dots \lambda_\chi$
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Singular value decomposition & compression

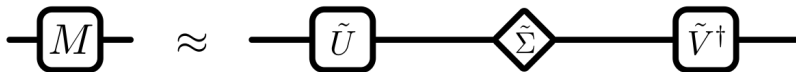


$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} = \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \begin{pmatrix} \lambda_1 & & & \\ & \lambda_2 & & \\ & & \lambda_3 & \\ & & & \lambda_4 \end{pmatrix} \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}$$

The diagram shows the SVD of a 4x4 matrix. The diagonal matrix Σ contains singular values $\lambda_1, \lambda_2, \lambda_3, \lambda_4$. The values λ_3 and λ_4 are crossed out with red X's, indicating they are being discarded for compression.

- Fixed bond dimension: Keep $\lambda_1 \dots \lambda_\chi$
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Singular value decomposition & compression



$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \approx \begin{pmatrix} \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \end{pmatrix} \begin{pmatrix} \lambda_1 & & & \\ & \lambda_2 & & \\ & & & \\ & & & \end{pmatrix} \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}$$

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Singular value decomposition & compression

$$\boxed{M} \approx \boxed{\tilde{U}\tilde{\Sigma}} \boxed{\tilde{V}^\dagger}$$

$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \approx \begin{pmatrix} \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \end{pmatrix} \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}$$

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$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \approx \begin{pmatrix} \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \end{pmatrix} \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}$$

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Matrix product states

$$|\psi\rangle = \sum_{i_1, i_2, \dots, i_N \in \{0,1\}} C_{i_1, i_2, \dots, i_N} |i_1, i_2, \dots, i_N\rangle$$

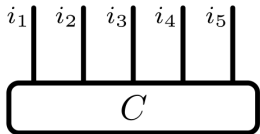
e.g. $C_{0,1,1,0,1} = \langle 01101 | \psi \rangle$

Reviews: Schollwöck, Ann. Phys. (N.Y.) 326, 96 (2011). Orús, *ibid.* 349, 117 (2014)

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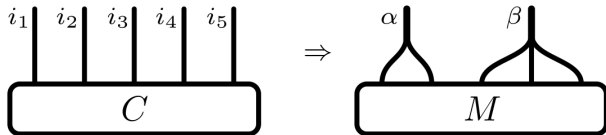


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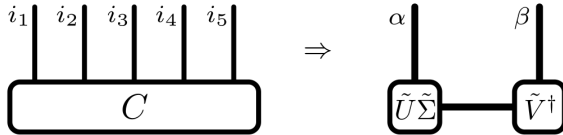


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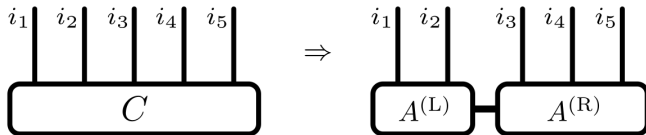


Reviews: Schollwöck, Ann. Phys. (N.Y.) 326, 96 (2011). Orús, *ibid.* 349, 117 (2014)

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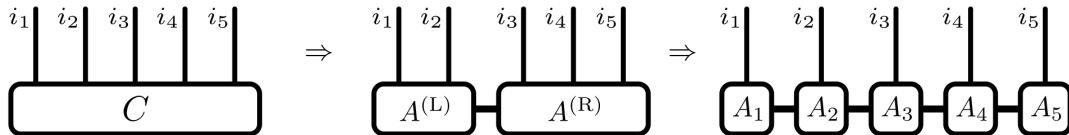


Reviews: Schollwöck, Ann. Phys. (N.Y.) 326, 96 (2011). Orús, *ibid.* 349, 117 (2014)

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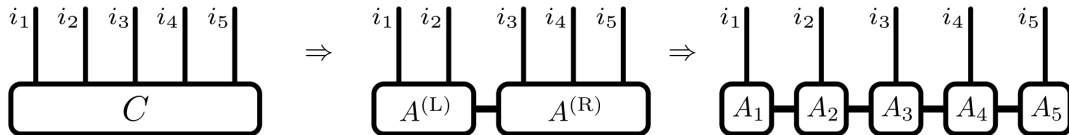


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Full state: 2^N numbers

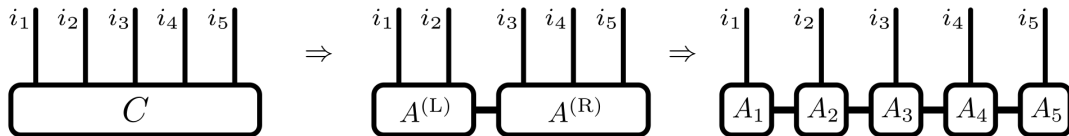
MPS, $2N\chi^2$ numbers

Reviews: Schollwöck, Ann. Phys. (N.Y.) 326, 96 (2011). Orús, *ibid.* 349, 117 (2014)

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Full state: 2^N numbers

MPS, $2N\chi^2$ numbers

Applies to: Wavefunctions, density matrices, probabilities

Also known as “tensor trains”

Reviews: Schollwöck, Ann. Phys. (N.Y.) 326, 96 (2011). Orús, *ibid.* 349, 117 (2014)

From Hilbert space to Liouville space

$$\hat{U}|\psi\rangle$$



From Hilbert space to Liouville space

$$\hat{U}|\psi\rangle\langle\psi|\hat{U}^\dagger$$



From Hilbert space to Liouville space

$$\hat{U}|\psi\rangle\langle\psi|\hat{U}^\dagger$$



$$\hat{U}\rho\hat{U}^\dagger$$



From Hilbert space to Liouville space

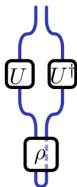
$$\hat{U}|\psi\rangle\langle\psi|\hat{U}^\dagger$$



$$\hat{U}\rho\hat{U}^\dagger$$



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From Hilbert space to Liouville space

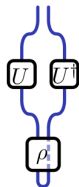
$$\hat{U}|\psi\rangle\langle\psi|\hat{U}^\dagger$$



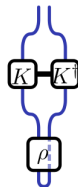
$$\hat{U}\rho\hat{U}^\dagger$$



$$\hat{U}\rho\hat{U}^\dagger$$



$$\sum_\alpha \hat{K}_\alpha \rho \hat{K}_\alpha^\dagger$$



From Hilbert space to Liouville space

Hilbert space

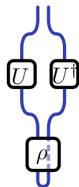
$$\hat{U}|\psi\rangle\langle\psi|\hat{U}^\dagger$$



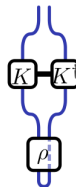
$$\hat{U}\rho\hat{U}^\dagger$$



$$\hat{U}\rho\hat{U}^\dagger$$



$$\sum_{\alpha} \hat{K}_{\alpha}\rho\hat{K}_{\alpha}^\dagger$$



From Hilbert space to Liouville space

Hilbert space

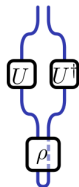
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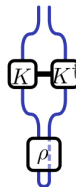
$$\hat{U}\rho\hat{U}^\dagger$$



$$\hat{U}\rho\hat{U}^\dagger$$



$$\sum_{\alpha} \hat{K}_{\alpha} \rho \hat{K}_{\alpha}^{\dagger}$$



Liouville space

$$\rightarrow \Lambda[\rho]$$

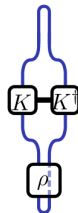
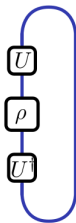
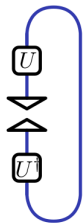
$\rightarrow$



From Hilbert space to Liouville space

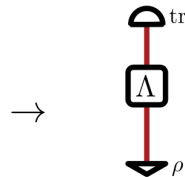
Hilbert space

$$\text{tr}(\hat{U}|\psi\rangle\langle\psi|\hat{U}^\dagger) \quad \text{tr}(\hat{U}\rho\hat{U}^\dagger) \quad \text{tr}(\hat{U}\rho\hat{U}^\dagger) \quad \text{tr}(\sum_\alpha \hat{K}_\alpha \rho \hat{K}_\alpha^\dagger)$$



Liouville space

$$\rightarrow \text{tr}(\Lambda[\rho])$$



QUAPI as tensor network

- So far ...: $\rho_{j_N}(t_N) = \sum_{j_1 \dots j_{N-1}} A^{j_N, j_{N-1}, \dots, j_1} \rho_{j_1}(t_1), \quad \mathbb{A} \rightarrow \mathbb{B} \cdot \mathbb{A}$

- B is tensor network: $B^{j_N, j_{N-1}, \dots, j_1}_{i_{N-1}, \dots, i_1} = [b_0]_{\alpha_1}^{j_N} \left(\prod_{k=1}^{n-2} [b_k]_{\alpha_{k+1}, i_{n-k}}^{\alpha_k, j_{n-k}} \right) [b_{n-1}]_{i_1}^{\alpha_{n-1}, j_1}$
- Finite memory approx — neglect correlations K steps ago

[Strathearn *et al.* Nat. Commun. '18]

QUAPI as tensor network

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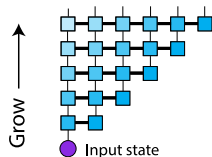
• Finite memory approx $\rightarrow$ neglect correlations K steps ago

[Strathearn *et al.* Nat. Commun. '18]

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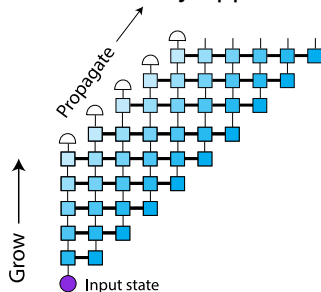
• Finite memory approx $\rightarrow$ neglect correlations K steps ago



[Strathearn *et al.* Nat. Commun. '18]

QUAPI as tensor network

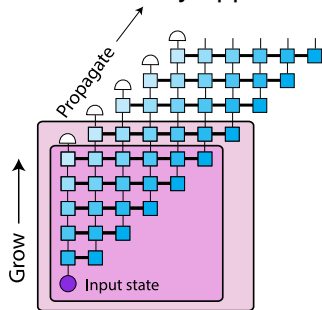
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[Strathearn *et al.* Nat. Commun. '18]

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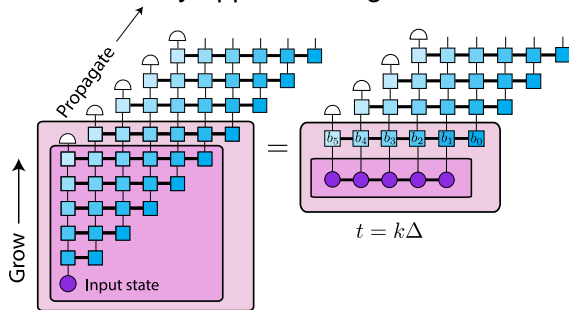
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[Strathearn *et al.* Nat. Commun. '18]

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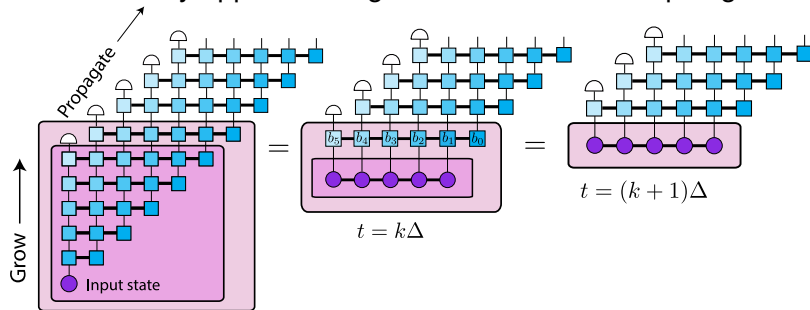
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[Strathearn *et al.* Nat. Commun. '18]

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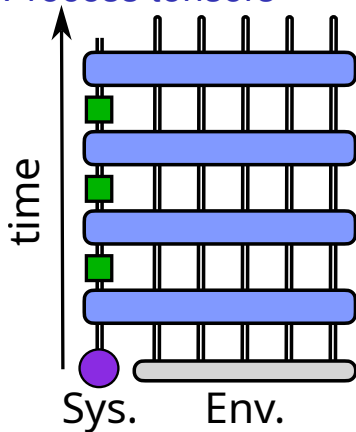
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[Strathearn *et al.* Nat. Commun. '18]

- 1 Non-Markovian problems & Many-body physics
 - Non-Markovian methods
- 2 Background: QUAPI & Tensor networks
 - Path integral approach
 - Tensor networks & Matrix product states
 - The TEMPO algorithm
- 3 Re-interpretation: Process tensors
 - PT-MPO representation
 - Beyond Gaussian environments
- 4 Applications
 - Optimisation
 - Lattice models
 - Polariton condensation
 - Two-dimensional spectroscopy
- 5 Future outlook

Process tensors

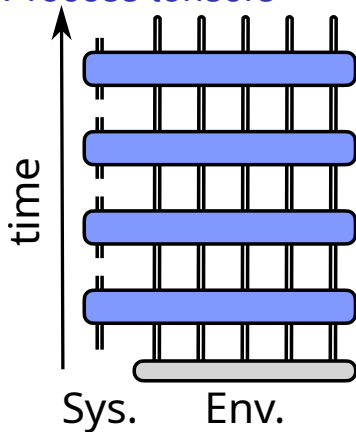


- Gate-based picture of evolution

- Separates out environment—response to sequence of operations
- Process tensor MPO

[Chiribella *et al.* PRL '08, PRA '09; Pollock *et al.* PRA '18, PRL '18]

Process tensors

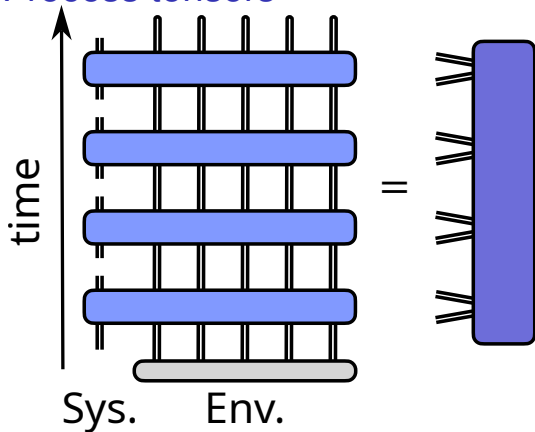


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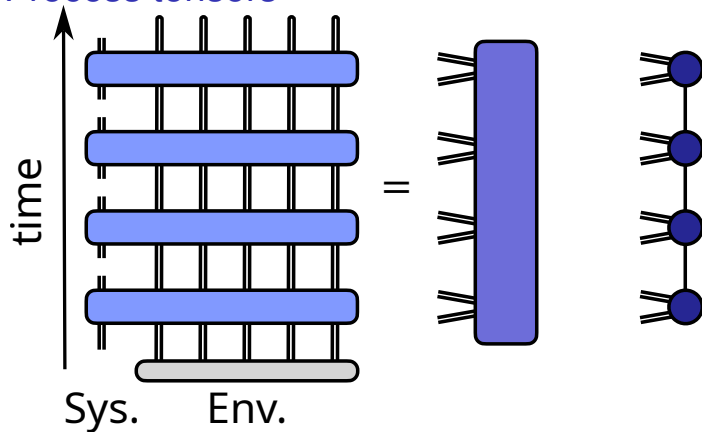


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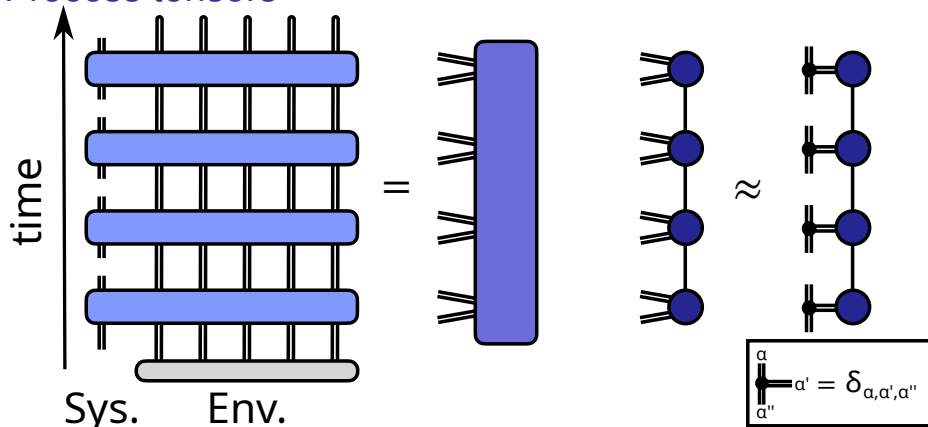
Process tensors



- Gate-based picture of evolution
- Separate out environment—response to sequence of operations
- Process tensor MPO

[Chiribella *et al.* PRL '08, PRA '09; Pollock *et al.* PRA '18, PRL '18]

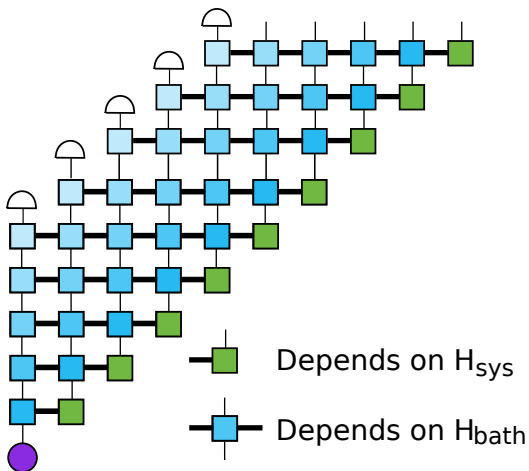
Process tensors



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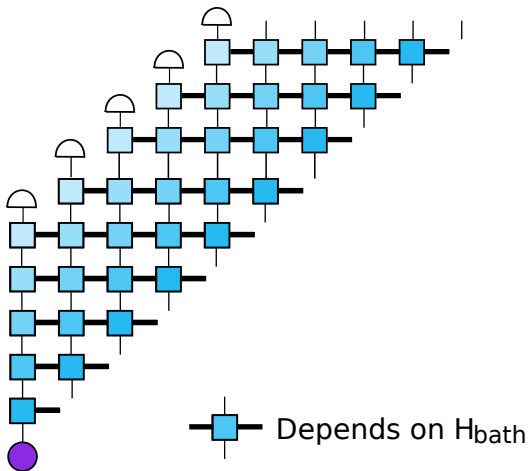
Re-interpreting TEMPO network $\rightarrow$ process tensors MPO



- Can separate bath/system terms

• Process tensor [Jørgensen & Pollock PRL '19]

Re-interpreting TEMPO network $\rightarrow$ process tensors MPO

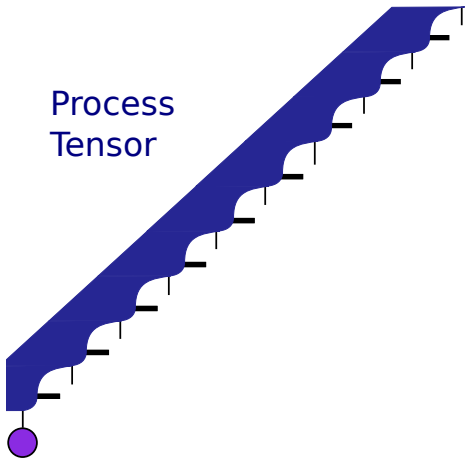


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Re-interpreting TEMPO network $\rightarrow$ process tensors MPO

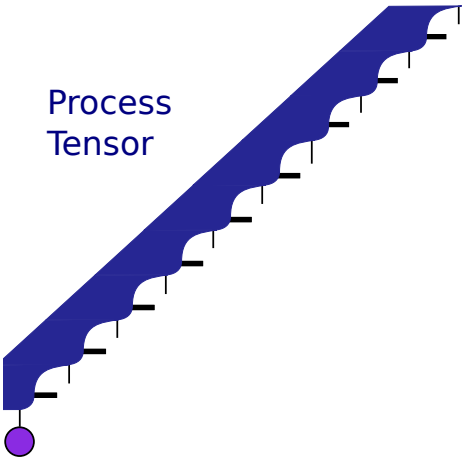
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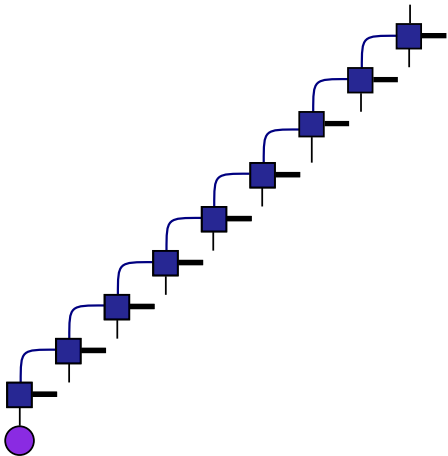
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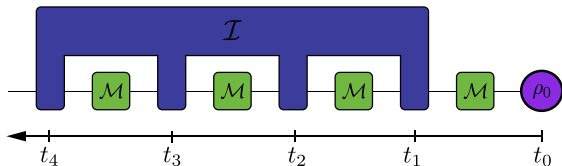
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Automatic compression of environments



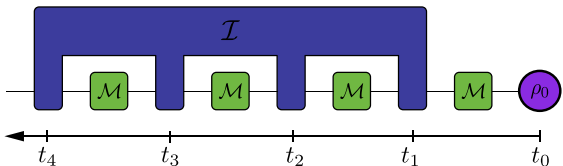
- Process tensor exists for any bath

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[Cygorek *et al.* Nat. Phys. '22]

Automatic compression of environments

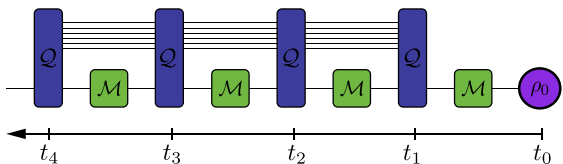


- Process tensor exists for any bath
- Explicit expression for PT exists
 - ▶ Sum over environment state space

▶ Compress to d via SVD

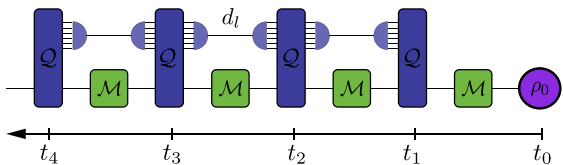
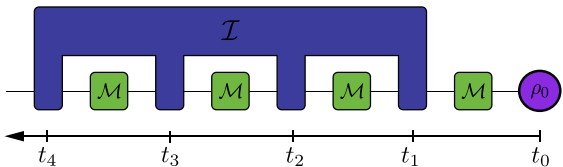
▶ Can build $Q[K]$ mode-by-mode

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[Cygorek *et al.* Nat. Phys. '22]

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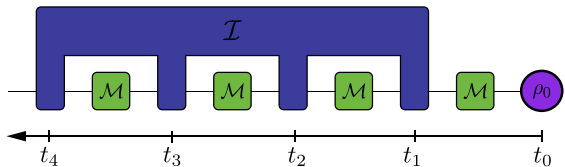
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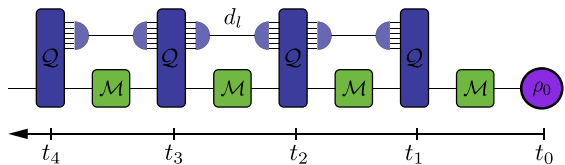
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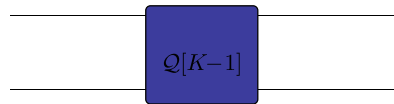
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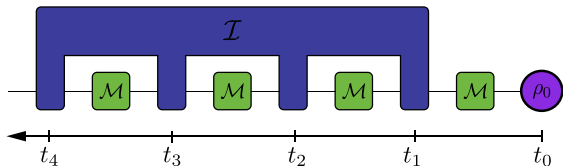


• Applications: Mode $\times$ crossover,
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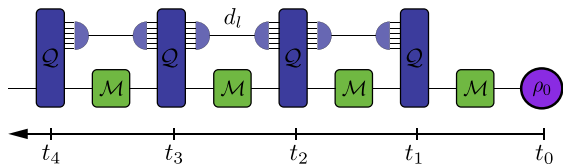


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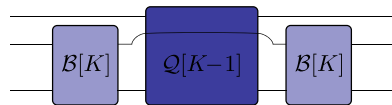
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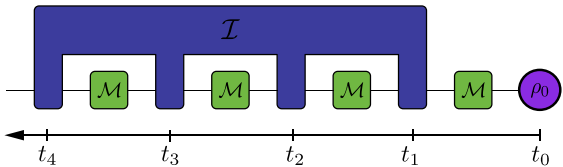


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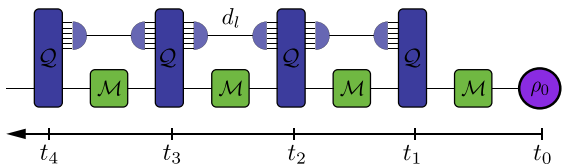


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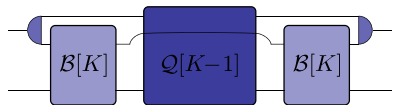
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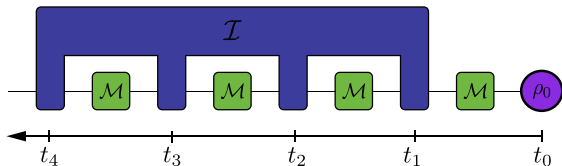


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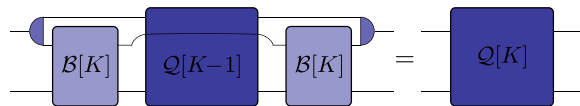
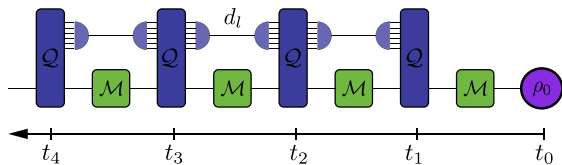
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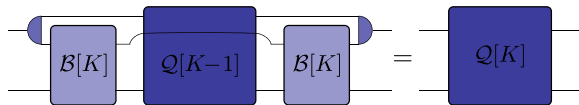
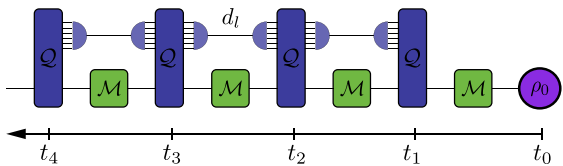
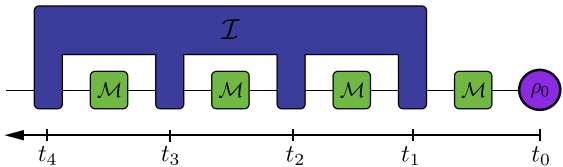
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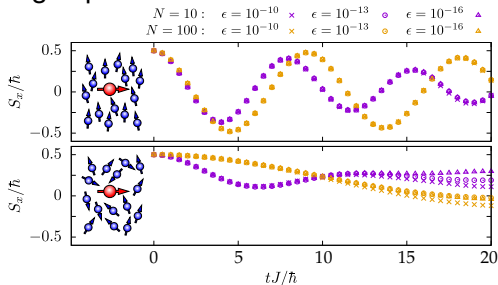
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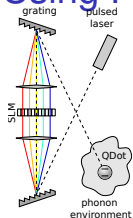
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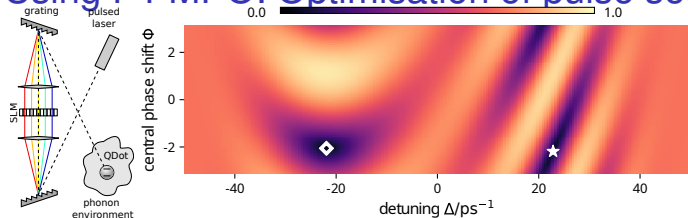
Using PT-MPO: Optimisation of pulse sequence



- 200s pre-compute, ~ 2 s contract $\times 16000$ points [Fux, Butler *et al.* PRL '21]
- Even faster: derivative w.r.t Hamiltonian via back-propagation

[Butler, Fux *et al.* PRL '24]

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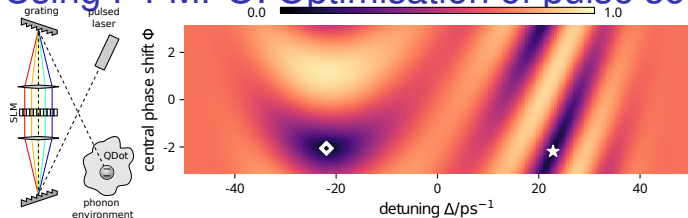


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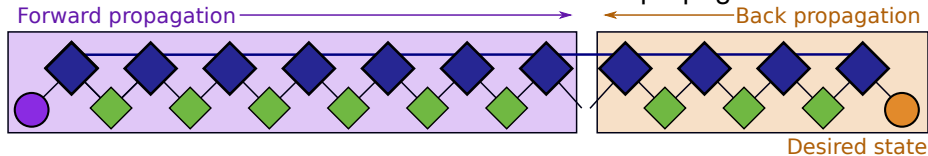
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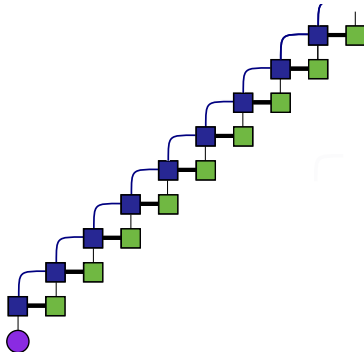
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TEMPO and Lattice problems

- One site: PT-MPO

• Chain of coupled sites

• e.g., thermalisation of 1D chain



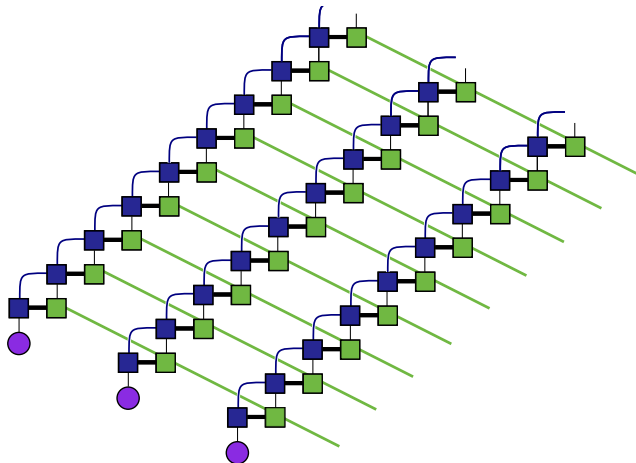
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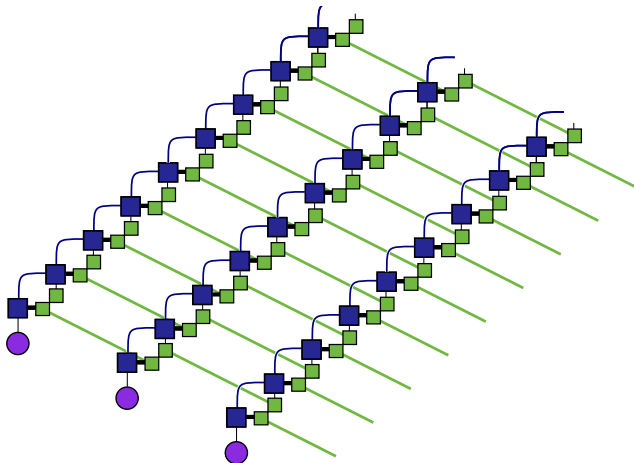
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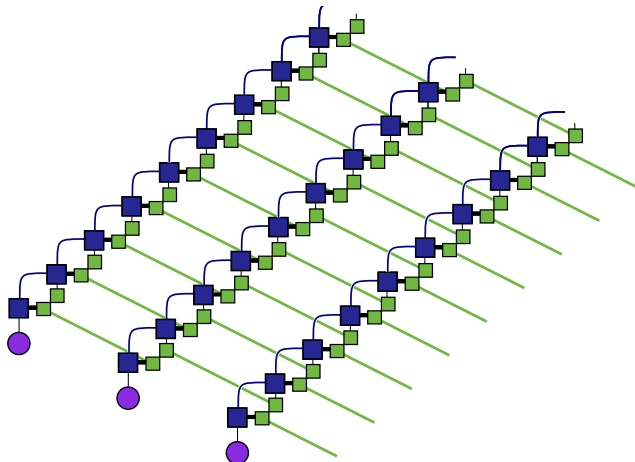
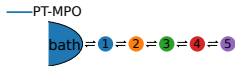
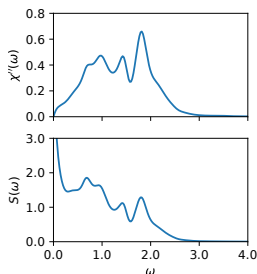
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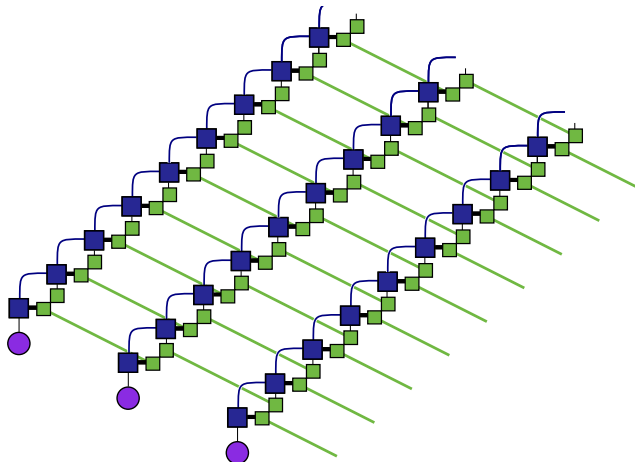
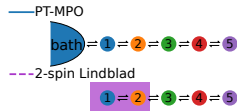
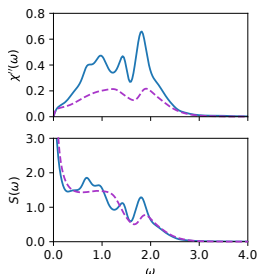
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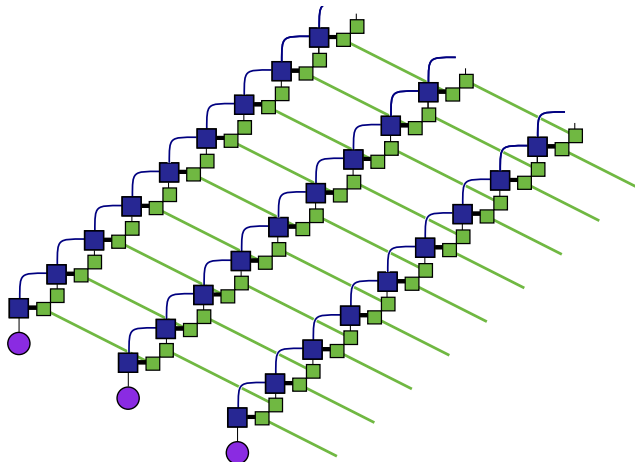
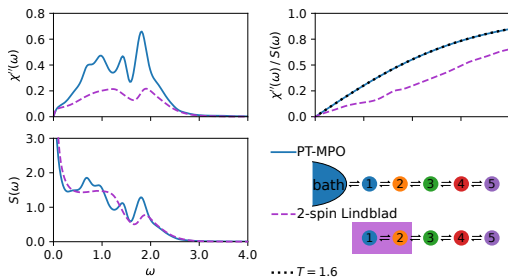
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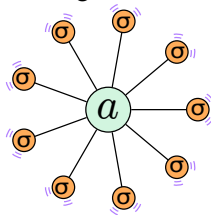


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Mean-field TEMPO: Polariton condensation

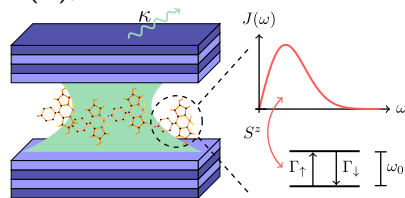
- Tavis-Cummings model + bath



$$H = \omega a^\dagger a + \sum_{i=1}^{N_m} \left[g(a\sigma_i^+ + \text{H.c.}) + \omega_0 \sigma_i^+ \sigma_i^- \right. \\ \left. + \sum_k \xi_{i,k} (b_{i,k} + b_{i,k}^\dagger) \sigma_i^z + \omega_{k,v} b_{i,k}^\dagger b_{i,k} \right]$$

- Large N_m , mean-field decoupling

- $J(\omega)$, match BODIPY-Br

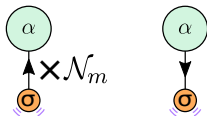


- Process tensor $\rightarrow$ fast evaluation
- Phase diagram vs $\omega_c, g, T_{\text{bath}}$

[Fowler-Wright *et al.* PRL '22] see also [Mu *et al.* PRL '26] beyond MF

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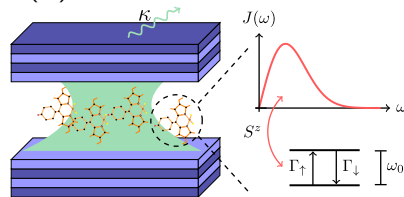


$$H_i = g(\langle a \rangle \sigma_i^x + \text{H.c.}) + \omega_0 \sigma_i^+ \sigma_i^-$$

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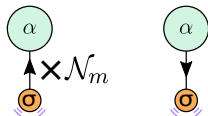


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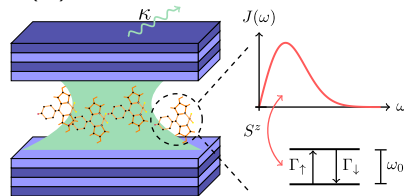


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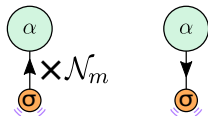


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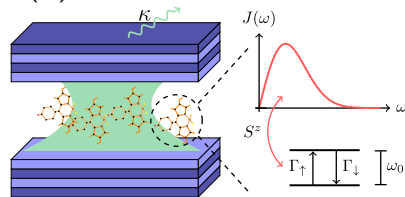


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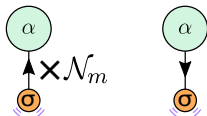
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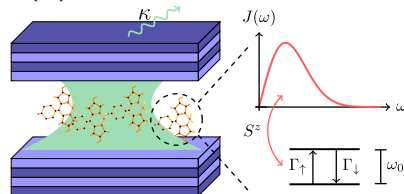


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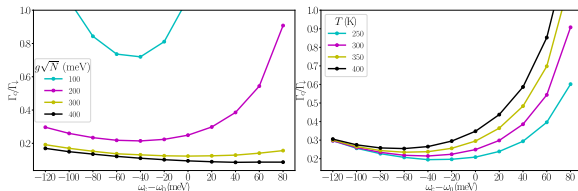
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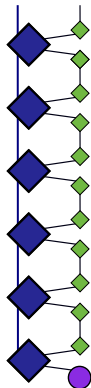
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Multi-time correlations $\rightarrow$ 2D spectroscopy

Process tensor:



• 2D spectroscopy.

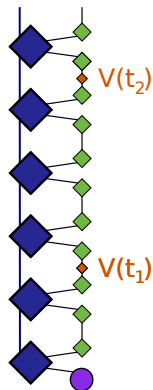
• 2D Spectra:

Fourier transform w.r.t $t_2 \rightarrow f_1$

[De Wit *et al.* PRR '25]

Multi-time correlations $\rightarrow$ 2D spectroscopy

Process tensor: $\rightarrow$ natural tool for multi-time correlations



[De Wit *et al.* PRR '25]

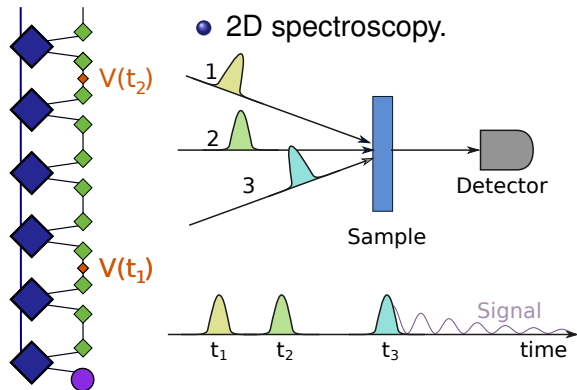
• 2D spectroscopy

• 2D Spectra:

Fourier transform w.r.t $t_2 \rightarrow f_2$

Multi-time correlations $\rightarrow$ 2D spectroscopy

Process tensor: $\rightarrow$ natural tool for multi-time correlations

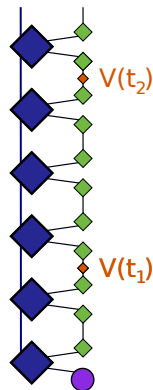


2D Spectra:
Fourier transform w.r.t $t_2 = t_1$

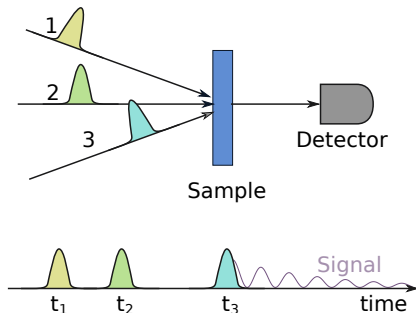
[De Wit *et al.* PRR '25]

Multi-time correlations $\rightarrow$ 2D spectroscopy

Process tensor: $\rightarrow$ natural tool for multi-time correlations



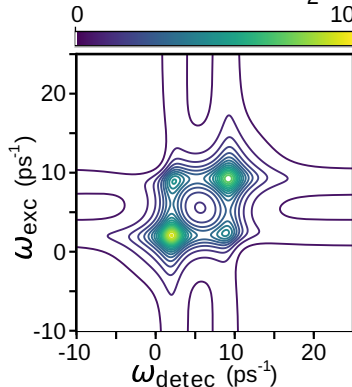
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[De Wit *et al.* PRR '25]

• 2D Spectra:

Fourier transform w.r.t $t_2 - t_1$



- 1 Non-Markovian problems & Many-body physics
 - Non-Markovian methods
- 2 Background: QUAPI & Tensor networks
 - Path integral approach
 - Tensor networks & Matrix product states
 - The TEMPO algorithm
- 3 Re-interpretation: Process tensors
 - PT-MPO representation
 - Beyond Gaussian environments
- 4 Applications
 - Optimisation
 - Lattice models
 - Polariton condensation
 - Two-dimensional spectroscopy
- 5 Future outlook

Outlook—future directions

- Complex systems: organic molecules, biological chromophores
Large Hilbert space, see [Cochin *et al.* 2603.06840]

• Many-body problems

• Process tensor representations beyond MPO

• Tree tensor network, e.g. [Dowling *et al.* PRX '24]

See [Keeling, Stoudenmire, Bañuls, Reichman: 2509.07661]

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 - ▶ Combine many body QQS & non-Markovian
 - ★ Effects on phase transitions
 - ★ Approaches: mean-field, lattice models, large Hilbert space
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 - ★ Combining environments, electrons + phonons
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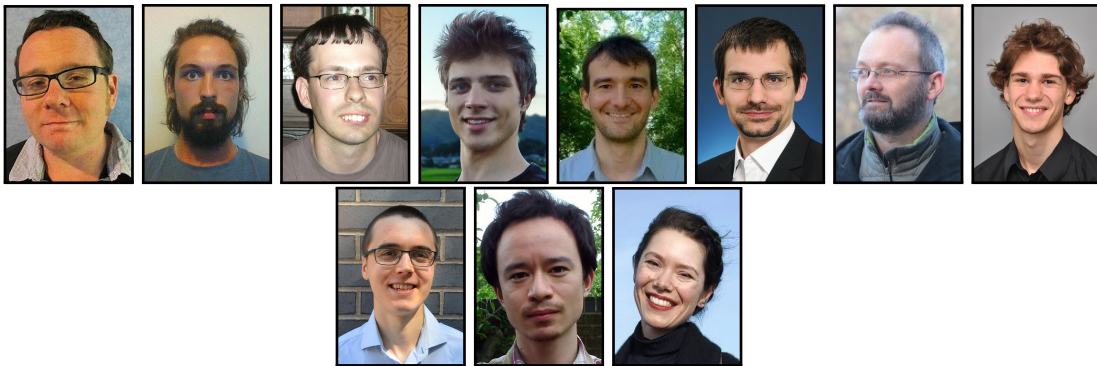
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Acknowledgements



FUNDING:

EPSRC

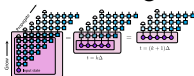
Engineering and Physical Sciences
Research Council



QOpen
Many-body Open Quantum Systems

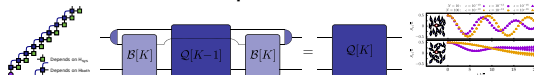
Summary

- TEMPO Algorithm for general non-Markovian problems



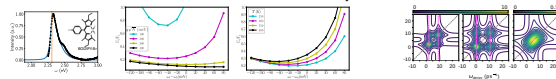
[Strathearn, Kirton, Kilda, Keeling & Lovett, Nat. Comm. '18]

- Reformulation as process tensor, Automatic Compression of Environments



[Fux *et al.* PRL '21] [Cygorek *et al.* Nat. Phys. '22]

- Can combine with Lattices, Mean-Field theory, 2D spectroscopy



[Fux, Kilda *et al.* PRR '23] [Fowler-Wright *et al.* PRL '22] [de Wit *et al.* PRR '25]

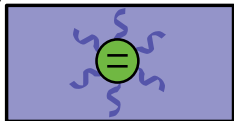
- Perspective: [Keeling *et al.* arXiv:2509.07661]
- **Open source code available:**

<https://github.com/tempoCollaboration/OQuPy> [Fux *et al.* JCP '24]

- 6 Application: Spin-Boson problem
 - Tensor networks & Matrix product states

spin Boson model

- Archetypal non-Markovian model:



$$H = \Omega S_x + \sum_i S_z (g_i a_i + g_i^* a_i^\dagger) + \omega_i a_i^\dagger a_i,$$

- Ohmic Bath density of states:

$$J(\omega) = \sum_i |g_i|^2 \delta(\omega - \omega_i) \equiv 2\alpha\omega \exp(-\omega/\omega_c)$$

- Known behaviour, initially excited [Leggett *et al.* RMP '87] for $\omega_c \gg \Omega$:

$0 < \alpha < 1/2$ Decaying oscillations, $\langle S_z \rangle \sim e^{i\omega t - \gamma t}$

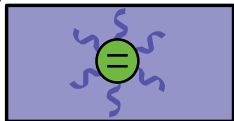
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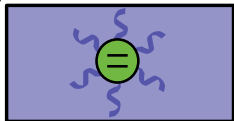
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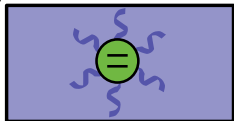
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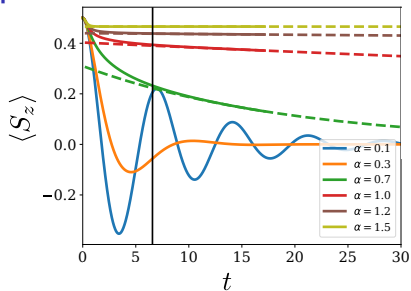
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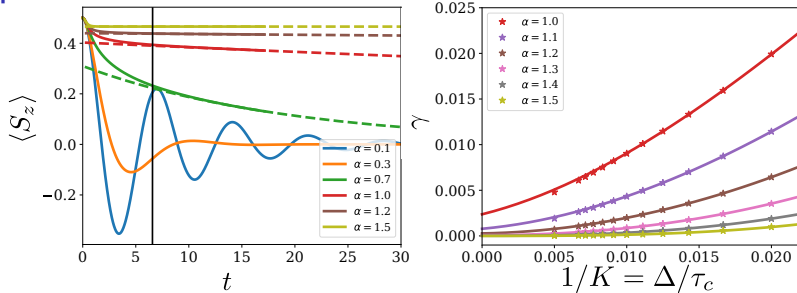
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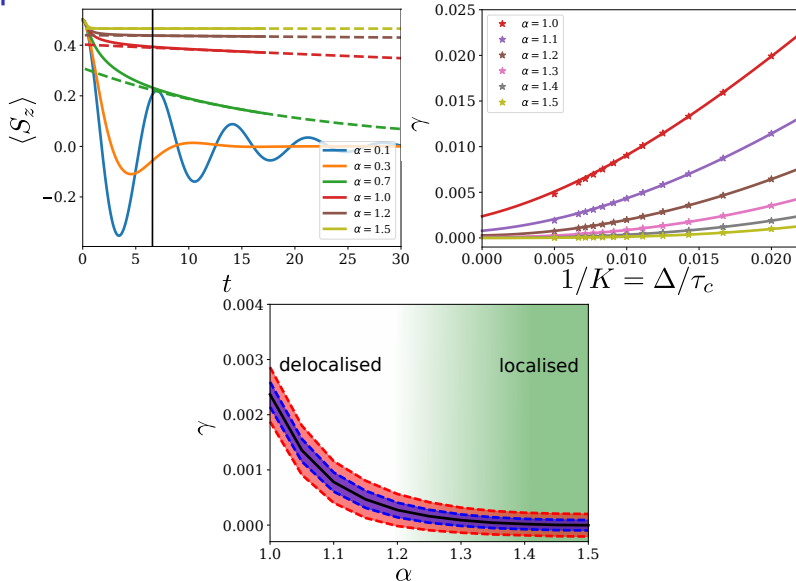
TEMPO spin Boson results



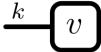
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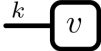
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Penrose tensor network notation & contraction

vector v_k 

Penrose tensor network notation & contraction

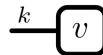
vector v_k A square box containing the letter 'v'. A horizontal line extends from the left side of the box, ending with the index 'k'.

matrix $A_{i,j}$ A square box containing the letter 'A'. Two horizontal lines extend from the box: one from the left side ending with the index 'i', and one from the right side ending with the index 'j'.

Penrose tensor network notation & contraction

vector

v_k



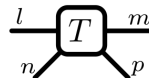
matrix

$A_{i,j}$



4-rank tensor

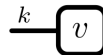
$T_{l,m,n,p}$



Penrose tensor network notation & contraction

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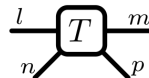
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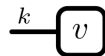
A



Penrose tensor network notation & contraction

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v_k



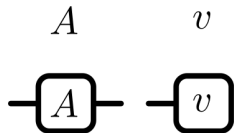
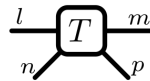
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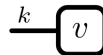


$A_{i,j}$ v_k

Penrose tensor network notation & contraction

vector

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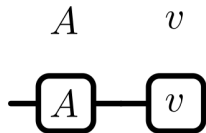
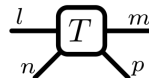
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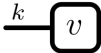
4-rank tensor

$T_{l,m,n,p}$

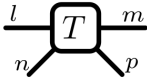


$$\sum_k A_{i,k} v_k$$

Penrose tensor network notation & contraction

vector v_k 

matrix $A_{i,j}$ 

4-rank tensor $T_{l,m,n,p}$ 

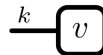
$$A \quad v = w$$


$$\sum_k A_{i,k} v_k = w_i$$

Penrose tensor network notation & contraction

vector

v_k



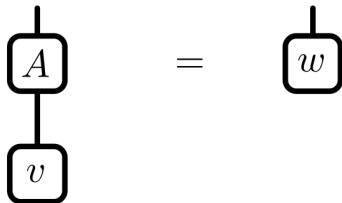
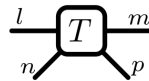
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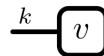
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Penrose tensor network notation & contraction

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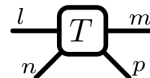
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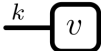


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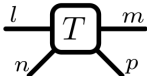
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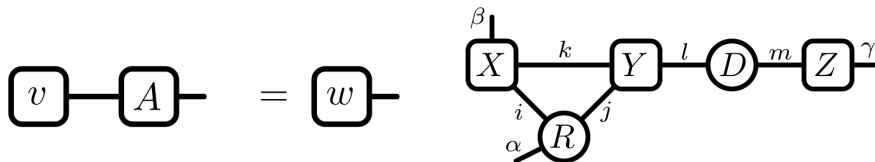


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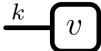
matrix $A_{i,j}$ 

4-rank tensor $T_{l,m,n,p}$ 

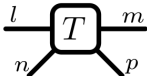


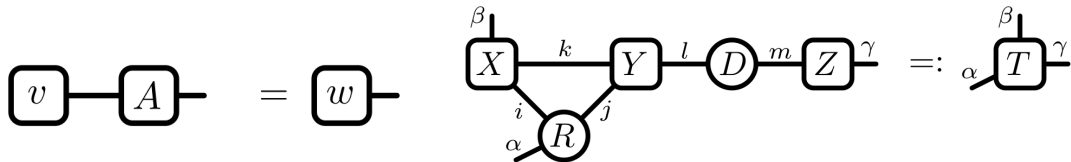
$$\sum_{ijklm} R_{\alpha ij} X_{i\beta k} Y_{jkl} D_{lm} Z_{m\gamma}$$

Penrose tensor network notation & contraction

vector v_k 

matrix $A_{i,j}$ 

4-rank tensor $T_{l,m,n,p}$ 



$$\sum_{ijklm} R_{\alpha ij} X_{i\beta k} Y_{jkl} D_{lm} Z_{m\gamma} =: T_{\alpha\beta\gamma}$$

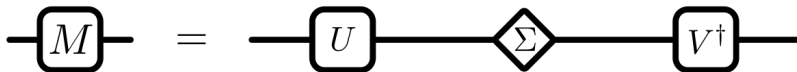
Singular value decomposition & compression



$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}$$

- Fixed bond dimension: Keep $\lambda_1, \dots, \lambda_\chi$
- Fixed precision: Keep $\lambda_n > \epsilon_{\text{rel}} \lambda_1$

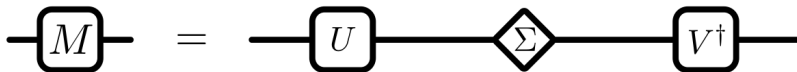
Singular value decomposition & compression



$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} = \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \begin{pmatrix} \lambda_1 & & & \\ & \lambda_2 & & \\ & & \lambda_3 & \\ & & & \lambda_4 \end{pmatrix} \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}$$

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Singular value decomposition & compression

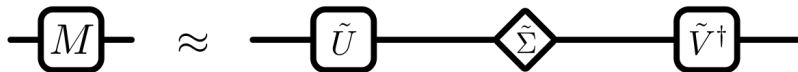


$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} = \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \begin{pmatrix} \lambda_1 & & & \\ & \lambda_2 & & \\ & & \lambda_3 & \\ & & & \lambda_4 \end{pmatrix} \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}$$

(Note: λ_3 and λ_4 in the diagonal matrix are crossed out with red X's)

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Singular value decomposition & compression



$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \approx \begin{pmatrix} \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \end{pmatrix} \begin{pmatrix} \lambda_1 & & & \\ & \lambda_2 & & \\ & & & \\ & & & \end{pmatrix} \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}$$

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Singular value decomposition & compression

$$\boxed{M} \approx \boxed{\tilde{U}\tilde{\Sigma}} \boxed{\tilde{V}^\dagger}$$

$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \approx \begin{pmatrix} \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \end{pmatrix} \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}$$

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Singular value decomposition & compression

$$\boxed{M} \approx \boxed{\tilde{U}\tilde{\Sigma}} \boxed{\tilde{V}^\dagger}$$

$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \approx \begin{pmatrix} \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \end{pmatrix} \begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}$$

- Fixed bond dimension: Keep $\lambda_1 \dots \lambda_\chi$

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Matrix product states

$$|\psi\rangle = \sum_{i_1, i_2, \dots, i_N \in \{0,1\}} C_{i_1, i_2, \dots, i_N} |i_1, i_2, \dots, i_N\rangle$$

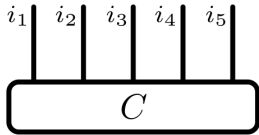
e.g. $C_{0,1,1,0,1} = \langle 01101 | \psi \rangle$

Reviews: Schollwöck, Ann. Phys. (N.Y.) 326, 96 (2011). Orús, *ibid.* 349, 117 (2014)

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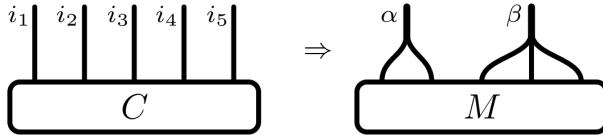


Reviews: Schollwöck, Ann. Phys. (N.Y.) 326, 96 (2011). Orús, *ibid.* 349, 117 (2014)

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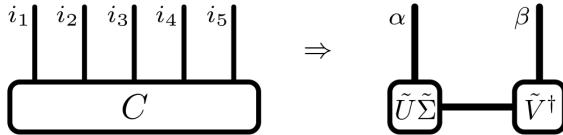


Reviews: Schollwöck, Ann. Phys. (N.Y.) 326, 96 (2011). Orús, *ibid.* 349, 117 (2014)

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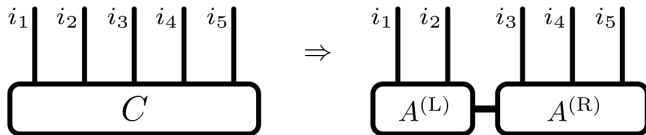


Reviews: Schollwöck, Ann. Phys. (N.Y.) 326, 96 (2011). Orús, *ibid.* 349, 117 (2014)

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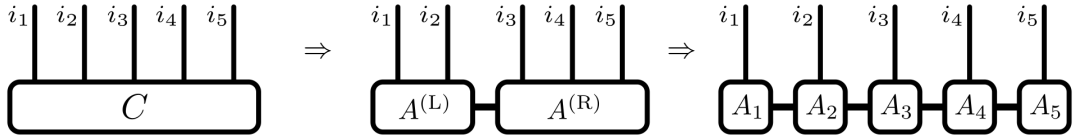


Reviews: Schollwöck, Ann. Phys. (N.Y.) 326, 96 (2011). Orús, *ibid.* 349, 117 (2014)

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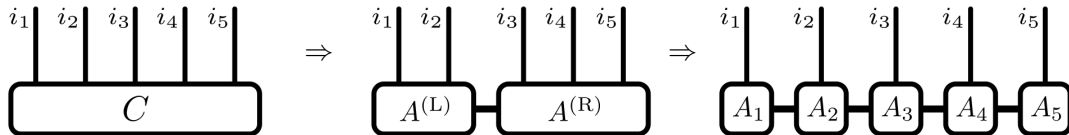


Reviews: Schollwöck, Ann. Phys. (N.Y.) 326, 96 (2011). Orús, *ibid.* 349, 117 (2014)

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Full state: 2^N numbers

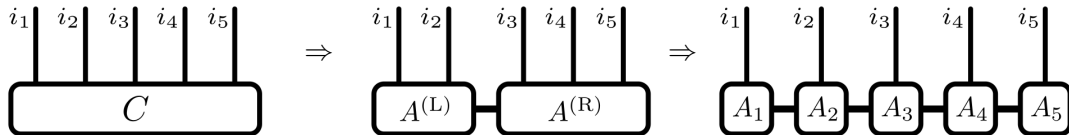
MPS, $2N\chi^2$ numbers

Reviews: Schollwöck, Ann. Phys. (N.Y.) 326, 96 (2011). Orús, *ibid.* 349, 117 (2014)

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Full state: 2^N numbers

MPS, $2N\chi^2$ numbers

Applies to: Wavefunctions, density matrices, probabilities

Also known as “tensor trains”

Reviews: Schollwöck, Ann. Phys. (N.Y.) 326, 96 (2011). Orús, *ibid.* 349, 117 (2014)

From Hilbert space to Liouville space

$$\hat{U}|\psi\rangle$$



From Hilbert space to Liouville space

$$\hat{U}|\psi\rangle\langle\psi|\hat{U}^\dagger$$



From Hilbert space to Liouville space

$$\hat{U}|\psi\rangle\langle\psi|\hat{U}^\dagger$$



$$\hat{U}\rho\hat{U}^\dagger$$



From Hilbert space to Liouville space

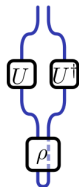
$$\hat{U}|\psi\rangle\langle\psi|\hat{U}^\dagger$$



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From Hilbert space to Liouville space

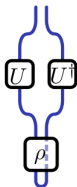
$$\hat{U}|\psi\rangle\langle\psi|\hat{U}^\dagger$$



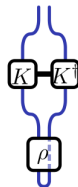
$$\hat{U}\rho\hat{U}^\dagger$$



$$\hat{U}\rho\hat{U}^\dagger$$



$$\sum_\alpha \hat{K}_\alpha \rho \hat{K}_\alpha^\dagger$$



From Hilbert space to Liouville space

Hilbert space

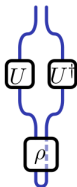
$$\hat{U}|\psi\rangle\langle\psi|\hat{U}^\dagger$$



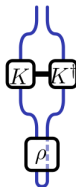
$$\hat{U}\rho\hat{U}^\dagger$$



$$\hat{U}\rho\hat{U}^\dagger$$



$$\sum_{\alpha} \hat{K}_{\alpha} \rho \hat{K}_{\alpha}^{\dagger}$$



From Hilbert space to Liouville space

Hilbert space

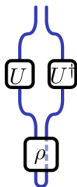
$$\hat{U}|\psi\rangle\langle\psi|\hat{U}^\dagger$$



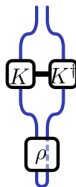
$$\hat{U}\rho\hat{U}^\dagger$$



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$$\sum_{\alpha} \hat{K}_{\alpha} \rho \hat{K}_{\alpha}^{\dagger}$$



Liouville space

$$\rightarrow \Lambda[\rho]$$

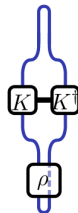
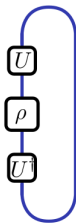
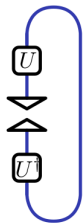
$\rightarrow$



From Hilbert space to Liouville space

Hilbert space

$$\text{tr}(\hat{U}|\psi\rangle\langle\psi|\hat{U}^\dagger) \quad \text{tr}(\hat{U}\rho\hat{U}^\dagger) \quad \text{tr}(\hat{U}\rho\hat{U}^\dagger) \quad \text{tr}(\sum_\alpha \hat{K}_\alpha \rho \hat{K}_\alpha^\dagger)$$



Liouville space

$$\rightarrow \text{tr}(\Lambda[\rho])$$

